

# Nonequilibrium Materials Engineering

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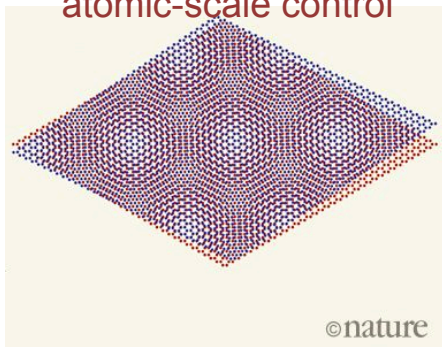
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Emmy Noether Programme (SE 2558/2-1)

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# Engineering materials with light

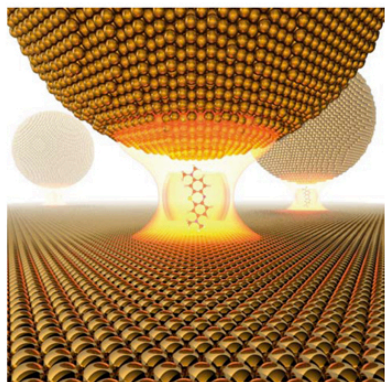
**condensed matter**  
quantum materials  
atomic-scale control



©nature

Y. Cao et al., Nature 556, 43 (2018)

**nonequilibrium materials engineering**



R. Chikkaraddy et al., Nature 535, 127 (2016)

**quantum optics**  
nanoplasmonics  
polaritonic chemistry

**QED: vacuum fluctuations**

**ultrafast spectroscopy**  
revealing elementary couplings  
light-induced new states of matter

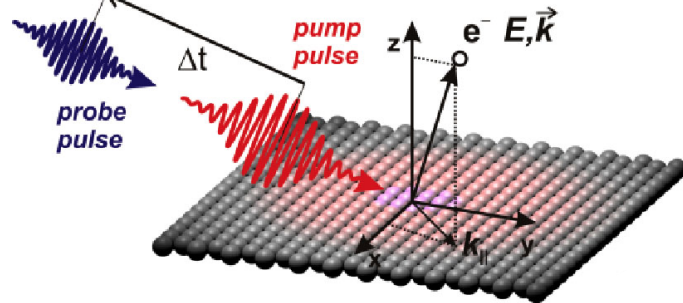
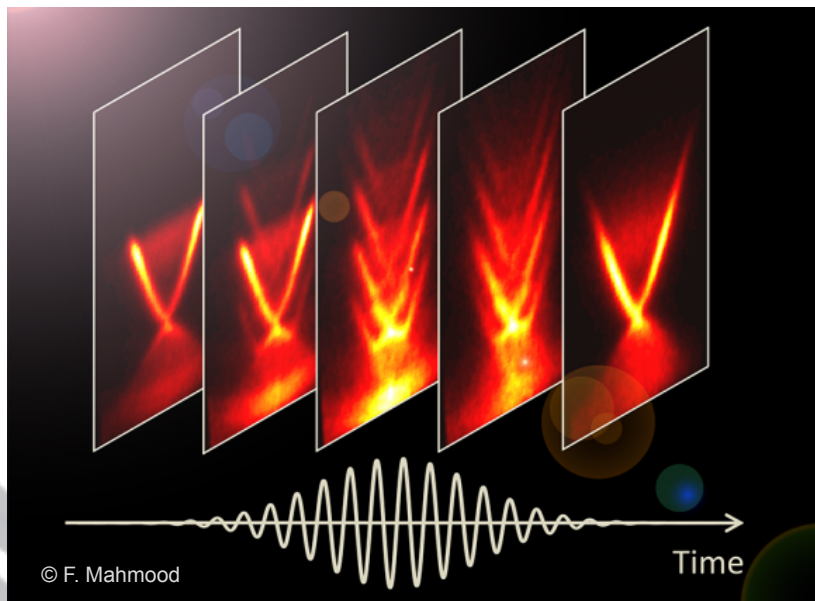


Image courtesy: J. Sobota

**pump-probe: strong classical fields**

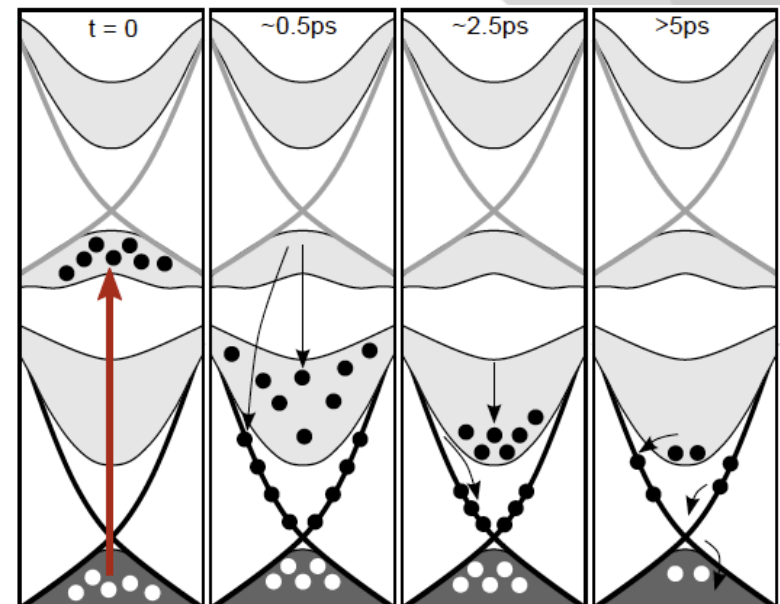
## Hamiltonian engineering

e.g., Floquet-Bloch bands



F. Mahmood et al., Nature Physics 12, 306 (2016)

## Distributional engineering



J. Sobota et al., JESRP 195, 249 (2014)

**many ingredients, hard to disentangle**

**this talk:** (I) tailored symmetry breaking, (II) vacuum fluctuations

## How to engineer materials with light?

### Part I: Optical control of chiral superconductors

Short laser pulses allow for switching of Majorana modes

*M. Claassen et al., Nat. Phys. 15, 766 (2019)*

### Part II: From classical to quantized photon fields

Materials engineering in an optical cavity using vacuum fluctuations

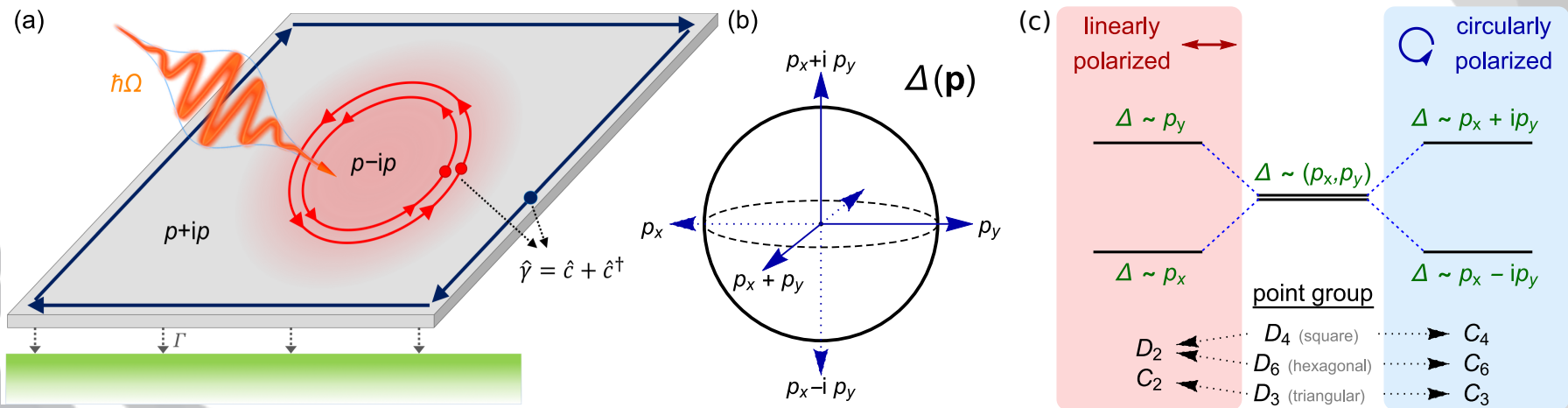
*M. A. Sentef et al., Science Advances 4, eaau6969 (2018)*



# I Optical control of Majoranas

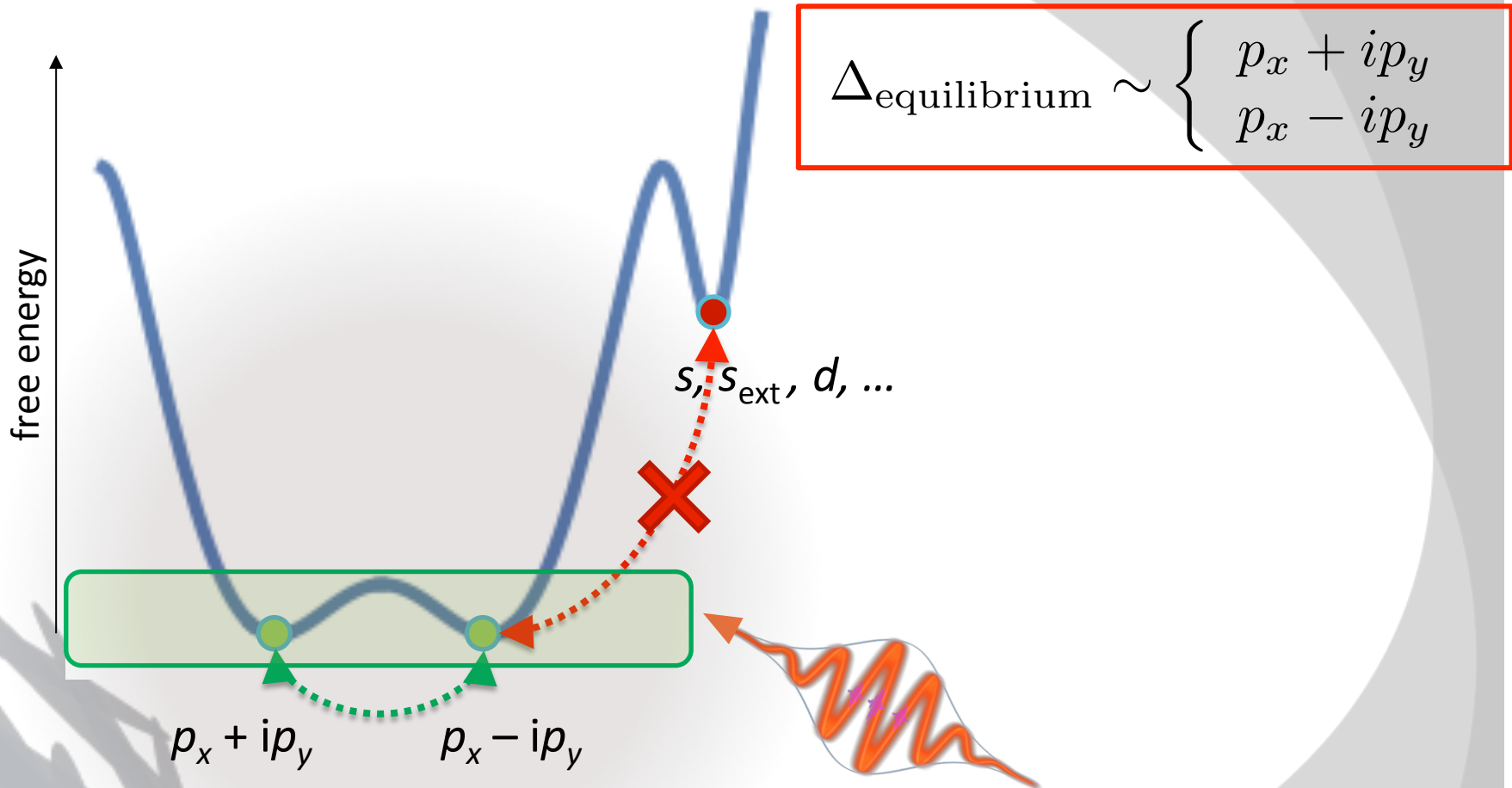
- can one switch the chirality of a 2D topological superconductor?

$\text{Sr}_2\text{RuO}_4(?)$ , highly doped graphene, twisted bilayer graphene, ...?



key idea: use two-pulse sequence with linearly and circularly polarized light

# Nonequilibrium pathway to switching



$$\Delta_{\text{non-eq}}(t) \sim \cos(\theta) \text{ "}p_x + ip_y\text{"} + \sin(\theta)e^{i\phi} \text{ "}p_x - ip_y\text{"}$$

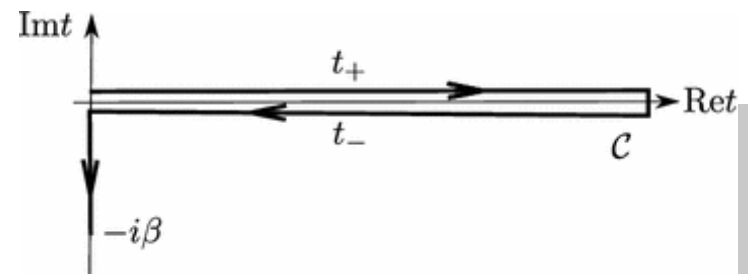
# Model and Method

multiband **Bogoliubov-de-Gennes** Hamiltonians for **doped graphene** (d+id) and **Sr2RuO4** (p+ip)  
coupling to **fermionic reservoir** to dissipate energy  
**laser driving** via Peierls substitution

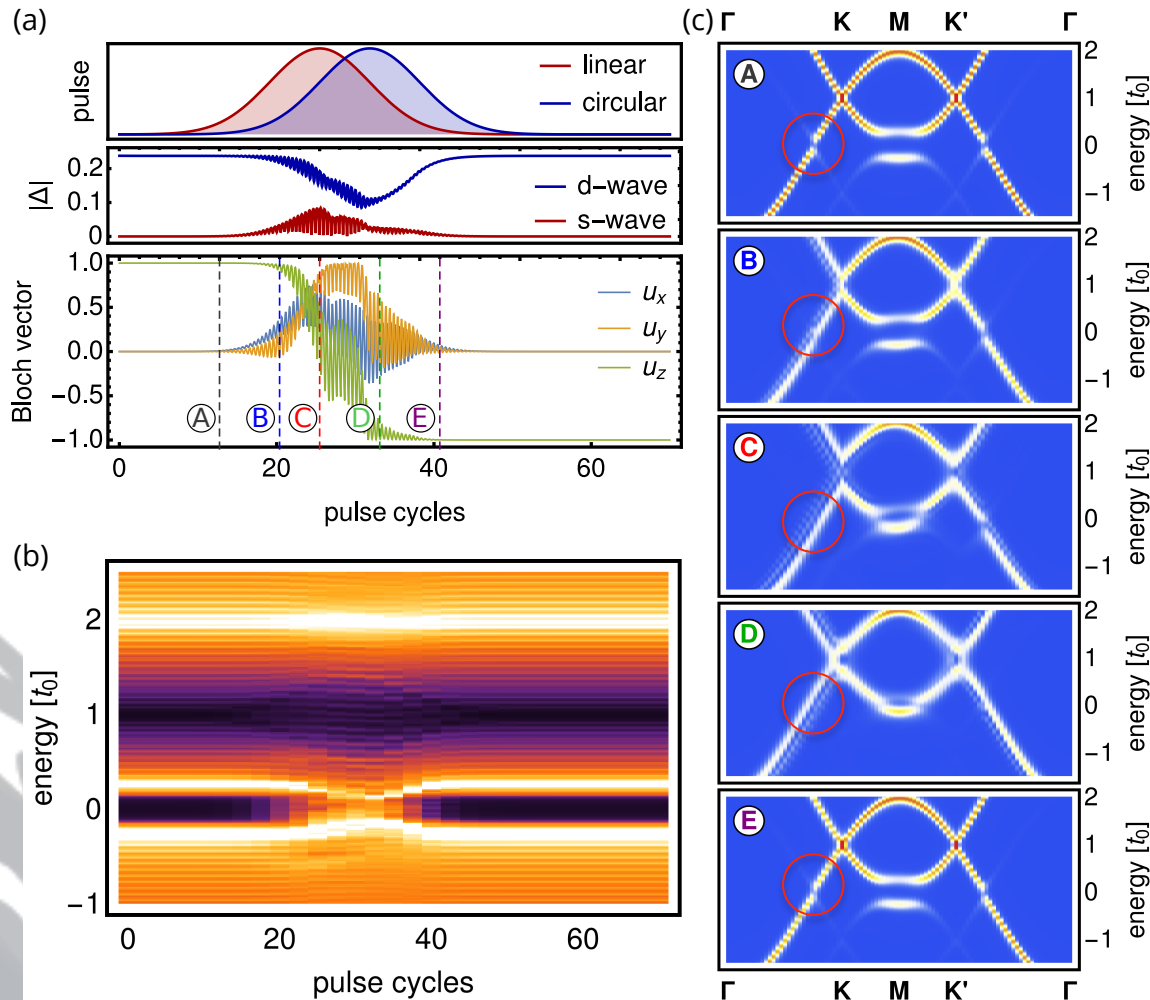
self-consistent Keldysh equations of motion for Nambu Green's functions:

$$i\partial_t \mathcal{G}_{\mathbf{k}}(t, t') = \mathcal{H}_{\mathbf{k}}(t, \Delta_{\mathbf{k}}(t)) \mathcal{G}_{\mathbf{k}}(t, t') + \int d\tau \hat{\Sigma}_{\mathbf{k}}(t, \tau) \mathcal{G}_{\mathbf{k}}(\tau, t')$$

$$\Delta_{\mathbf{k}}(t) = \frac{1}{L} \sum_j v^{(j)} \hat{\eta}_{\mathbf{k}}^{(j)} \sum_{\substack{\mathbf{k}' \\ \alpha\beta}} \hat{\eta}_{\mathbf{k}'\alpha\beta}^{(j)} \left\langle \hat{c}_{-\mathbf{k}',\beta\downarrow} \hat{c}_{\mathbf{k}',\alpha\uparrow} \right\rangle$$



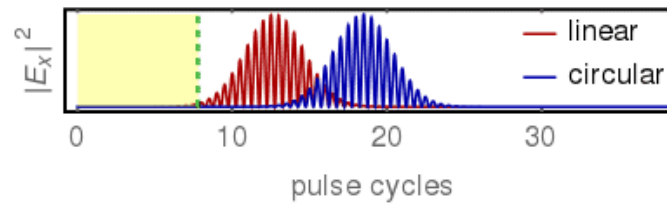
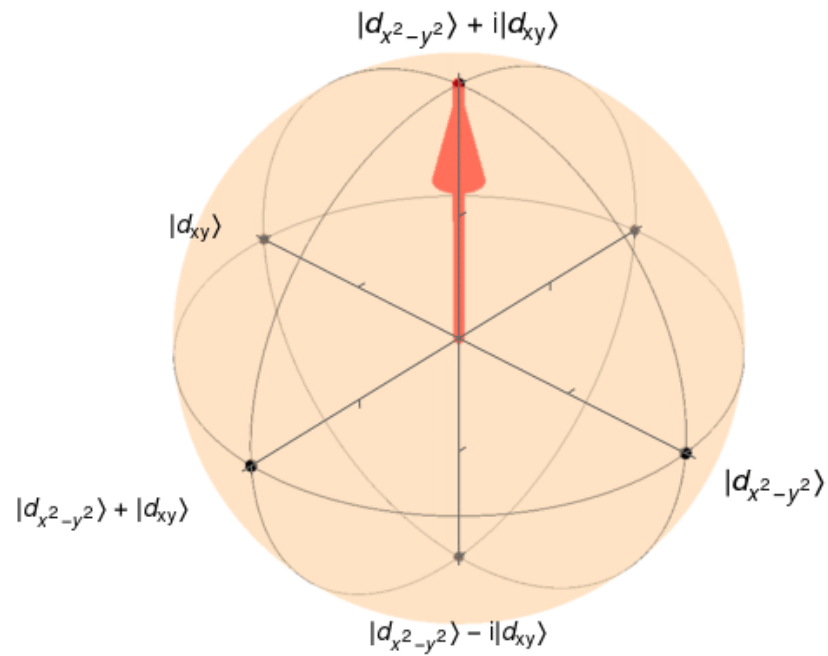
# Optical control of Majoranas



two-pulse sequence  
reverses d+id state  
in graphene

time-resolved  
spectroscopy tracks  
chirality reversal

# Bloch vector rotation

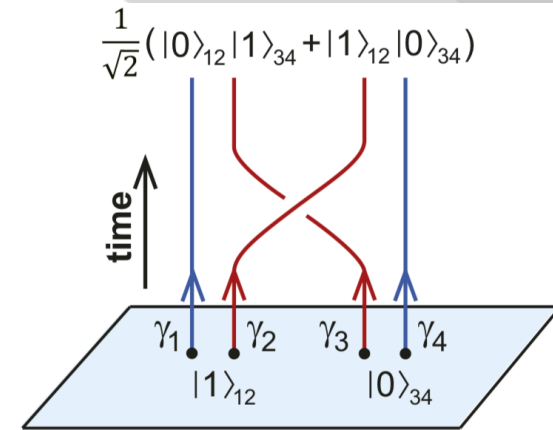




# A „programmable“ topological quantum computer?

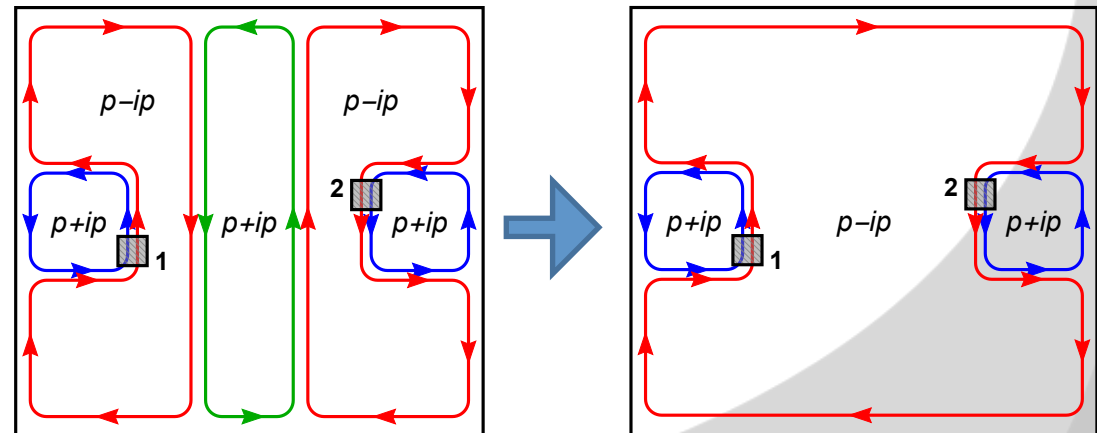
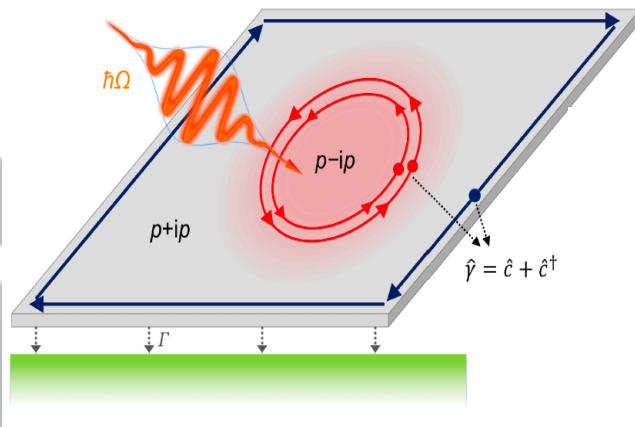
## non-Abelian statistics of Majorana fermions:

- half-quantum vortices of chiral superconductors host single Majorana fermions
  - Two Majoranas represent one electron:  $\frac{1}{2} + \frac{1}{2} = 1$
- Braiding between Majoranas is a non-Abelian operation in electron (charge) basis!



Ivanov, PRL 86, 268 (2001)  
B. Lian et al., PNAS 115, 10938 (2018)

simplest operation: a **switchable Hadamard gate**



# Summary I

- All-optical **control of chiral Majorana modes**
- towards arbitrarily programmable quantum computer?

„program the gate optically, read it out electrically“

*M. Claassen et al.,  
Nat. Phys. 15, 766 (2019)*



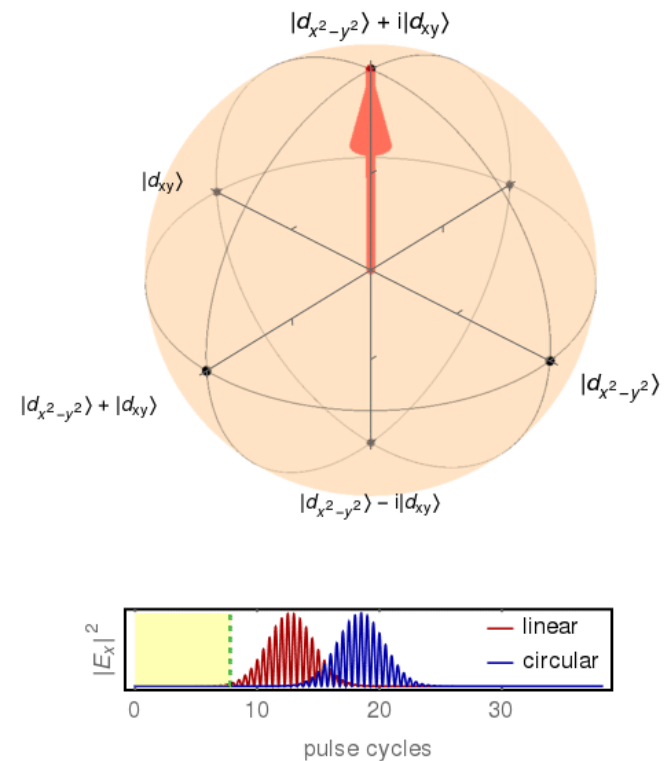
M. Claassen



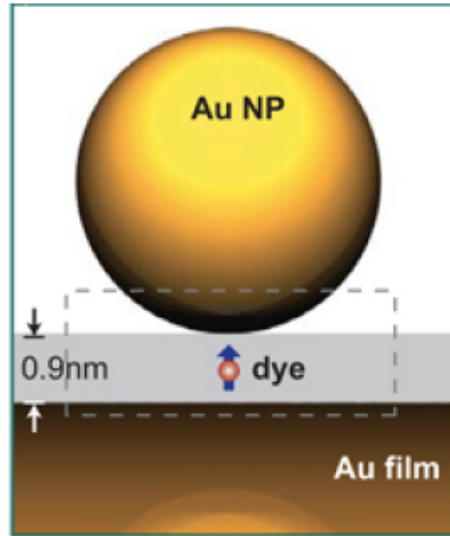
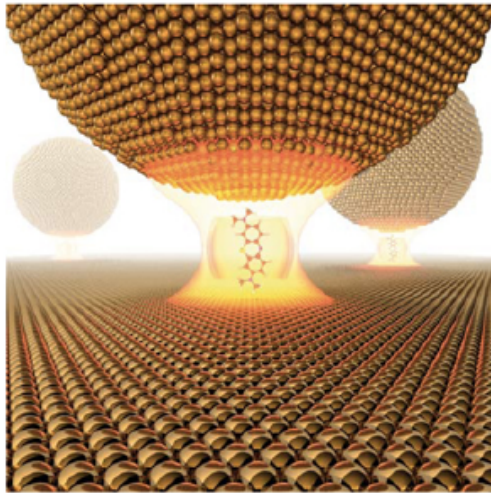
D. Kennes



M. Zingl



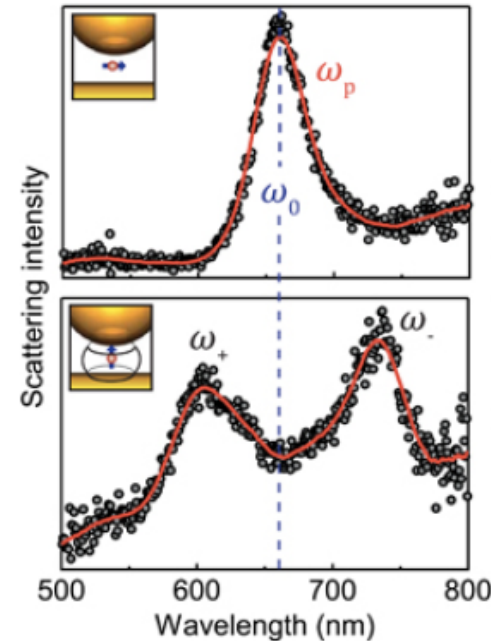
# From classical to quantum light



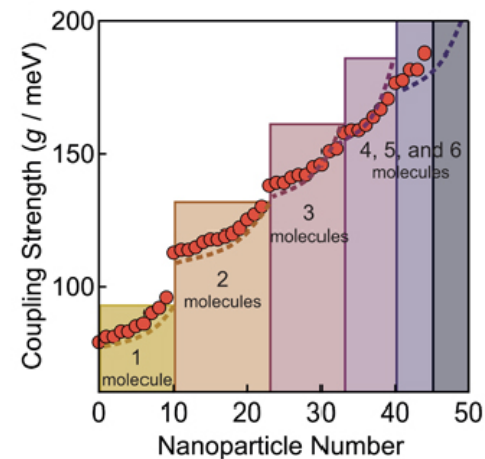
R. Chikkaraddy et al., Nature 535, 127 (2016)

**collective** strong light-matter coupling

what about **cavity materials**?

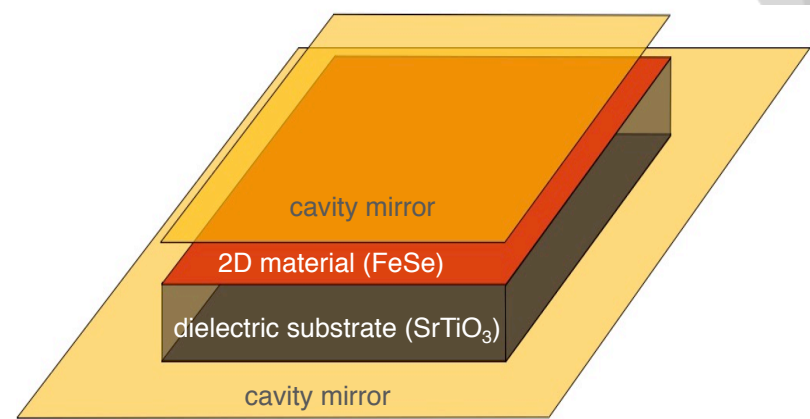
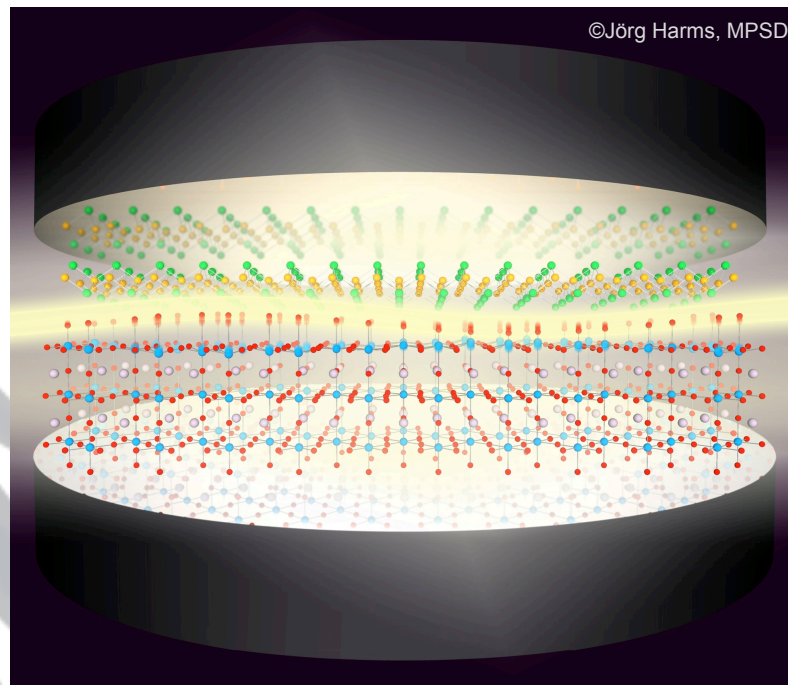


**Rabi splitting**

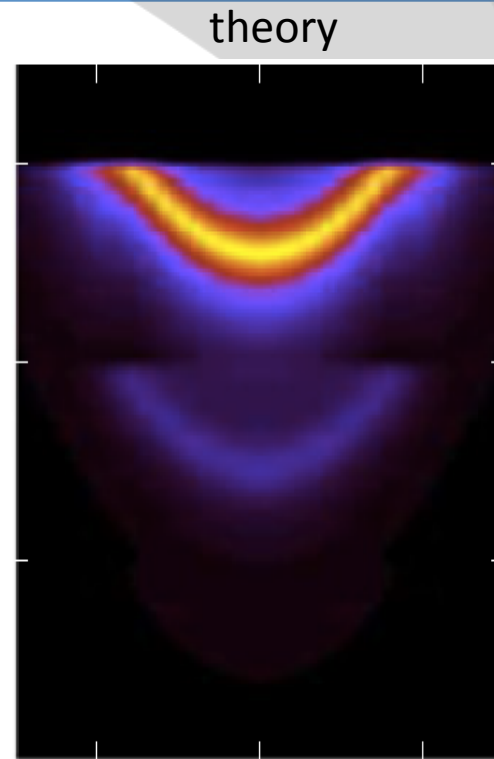
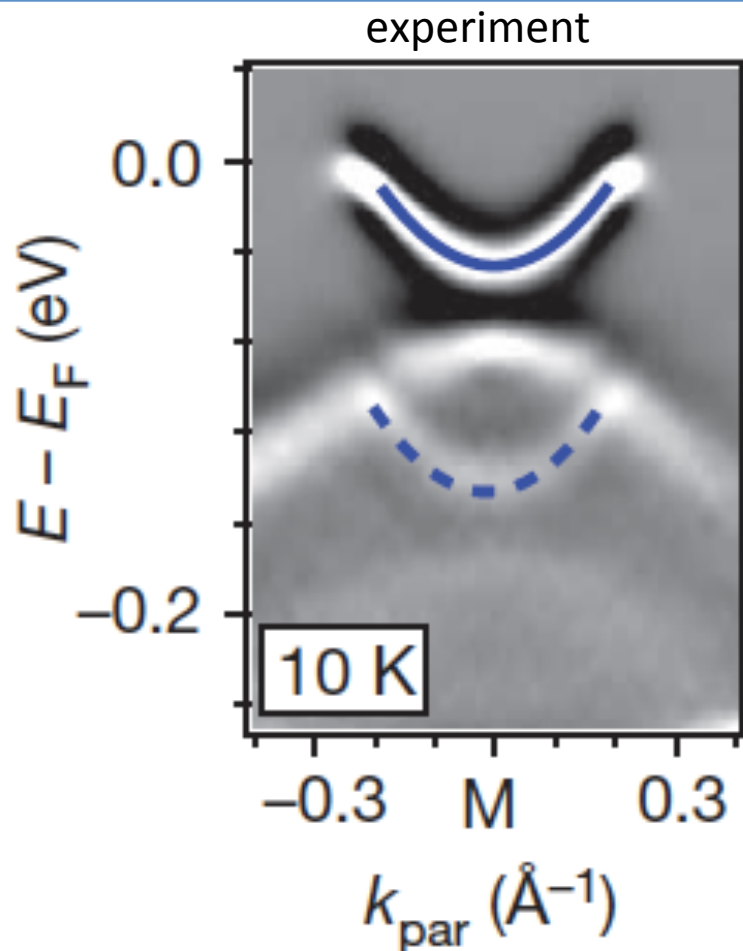


# II Cavity materials

- can one use enhanced vacuum fluctuations to change materials properties?



# monolayer FeSe/STO: ARPES



replica bands: forward (small- $q$ )  
electron-phonon scattering

*Lee et al., Nature 515, 245 (2014)*

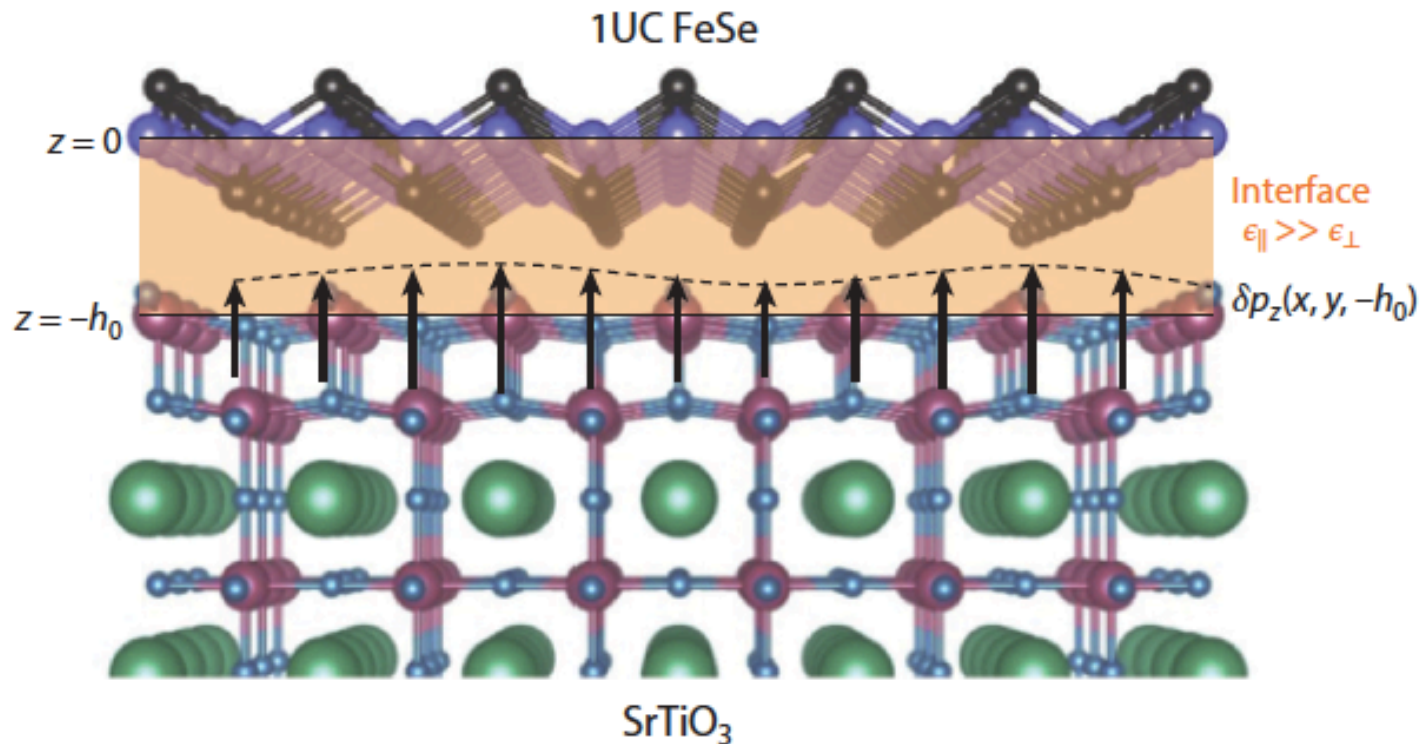
*Rademaker et al., New J. Phys. 18, 022001 (2016)*



# monolayer FeSe/STO: interfacial phonon

bare el-phonon vertex  $g(\vec{q}) = g_0 \exp(-|\vec{q}|/q_0)$  *Lee et al., Nature 515, 245 (2014)*

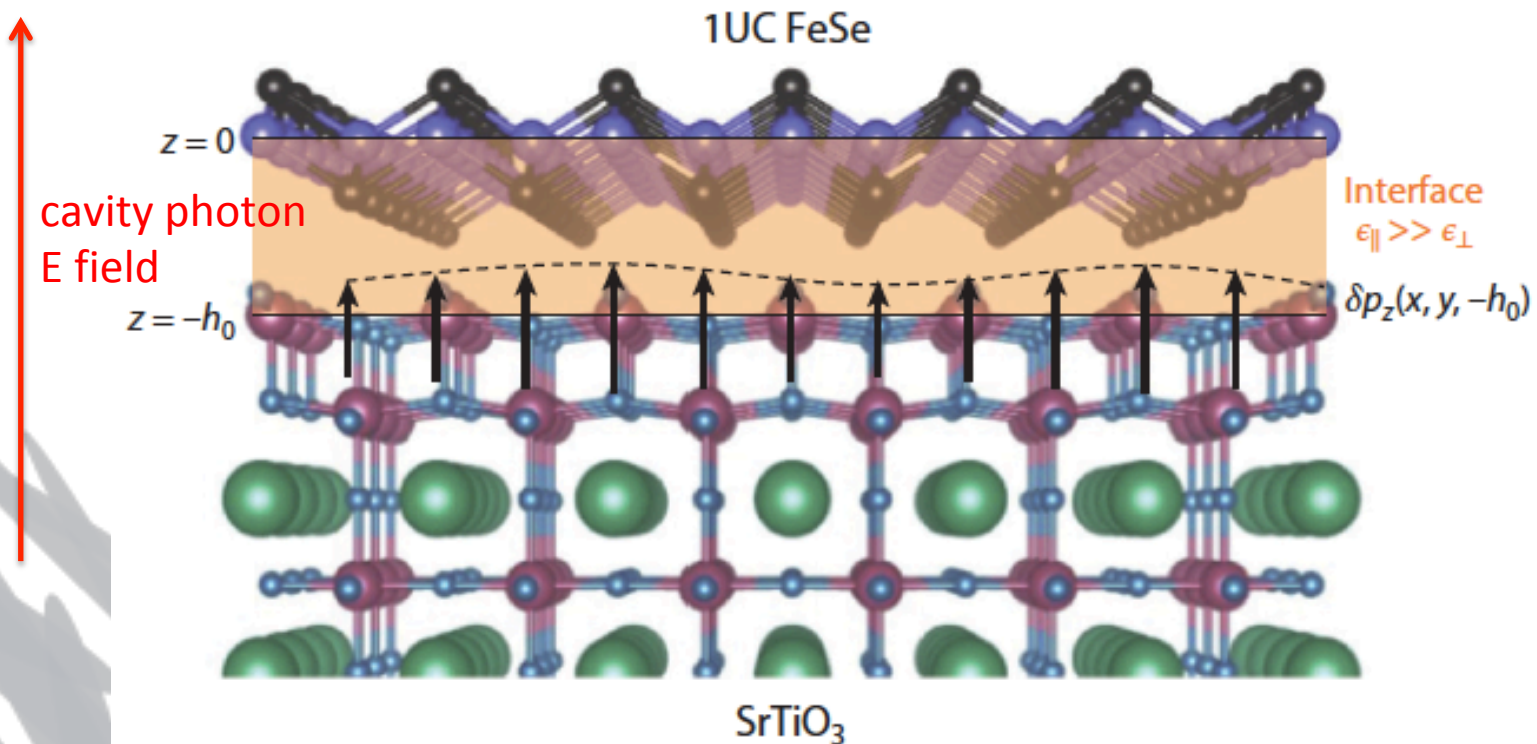
$$q_0^{-1} = h_0 \sqrt{\epsilon_{\parallel}/\epsilon_{\perp}} \quad \epsilon_{\parallel}/\epsilon_{\perp} \approx 100$$



*Huang and Hoffman, Annu. Rev. CMP 8, 311 (2017)*

# Cavity engineering

- idea: use **phonon polaritons** to enhance electron-phonon coupling



*Huang and Hoffman, Annu. Rev. CMP 8, 311 (2017)*

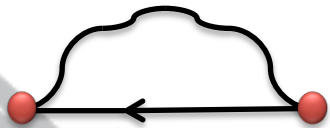
$$H = \sum_{\vec{k}, \sigma} \epsilon_{\vec{k}} c_{\vec{k}, \sigma}^\dagger c_{\vec{k}, \sigma} + \frac{1}{\sqrt{N}} \sum_{\vec{k}, \vec{q}, \sigma, \lambda = \pm} c_{\vec{k} + \vec{q}, \sigma}^\dagger c_{\vec{k}, \sigma} (g_\lambda^*(\vec{q}) \alpha_{-\vec{q}, \lambda}^\dagger + g_\lambda(\vec{q}) \alpha_{\vec{q}, \lambda}) + \sum_{\vec{q}, \lambda = \pm} \omega_\lambda(\vec{q}) \alpha_{\vec{q}, \lambda}^\dagger \alpha_{\vec{q}, \lambda}$$

electrons                      el-polariton coupling                      polaritons

bare el-phonon vertex     $g(\vec{q}) = g_0 \exp(-|\vec{q}|/q_0)$      $q_0^{-1} = h_0 \sqrt{\epsilon_{\parallel} / \epsilon_{\perp}}$

**G-self-consistent** Migdal-Eliashberg theory

$$\hat{\Sigma}(\vec{k}, i\omega_n) = \frac{-1}{N\beta} \sum_{\vec{q}, m, \lambda = \pm} |g_\lambda(\vec{q})|^2 D_\lambda^{(0)}(\vec{q}, i\omega_n - i\omega_m) \hat{\tau}_3 \hat{G}(\vec{k} + \vec{q}, i\omega_m) \hat{\tau}_3$$

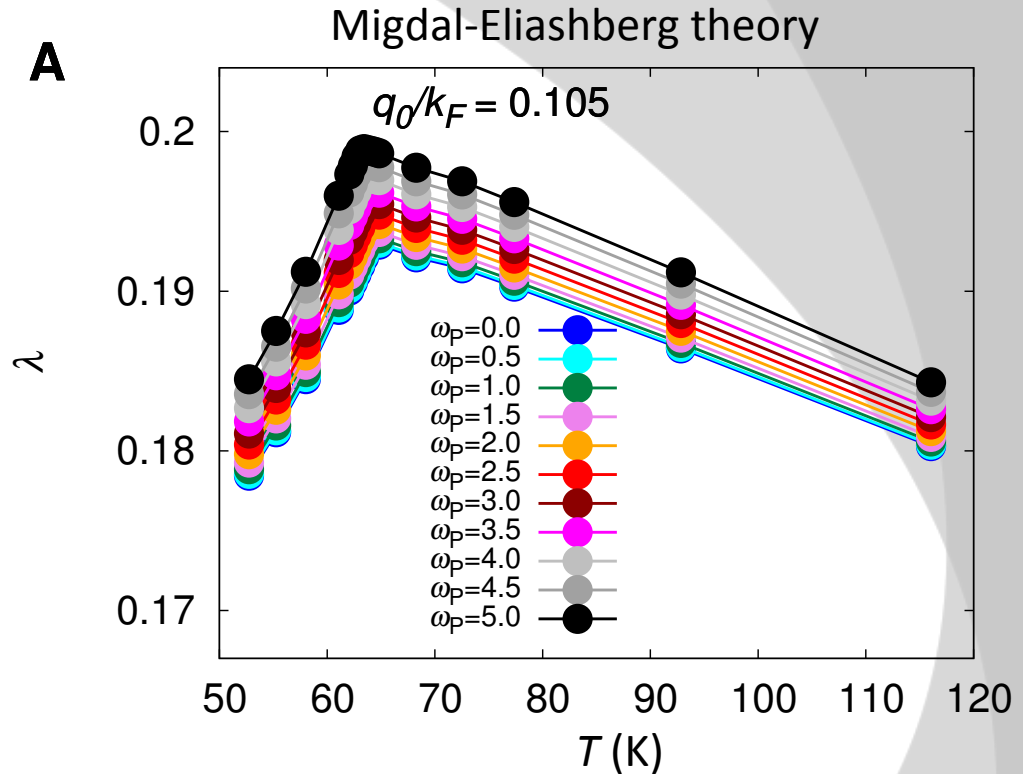
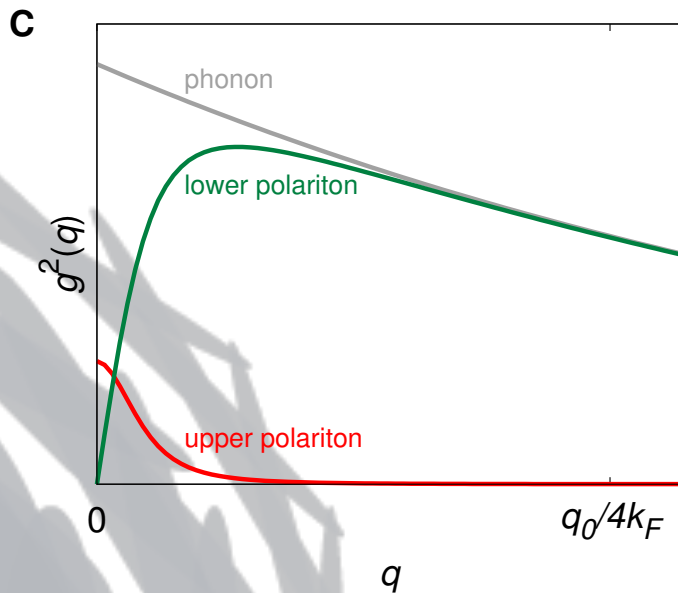
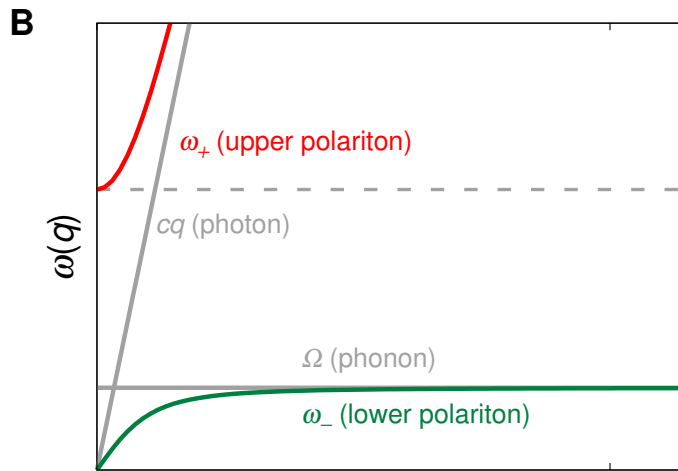


$$\hat{\Sigma}(\vec{k}, i\omega_n) = i\omega_n [1 - Z(\vec{k}, i\omega_n)] \hat{\tau}_0 + \chi(\vec{k}, i\omega_n) \hat{\tau}_3 + \phi(\vec{k}, i\omega_n) \hat{\tau}_1$$

$$\lambda \equiv Z(\vec{k}_F, i\pi/\beta) - 1$$

Mass enhancement:  $m^*/m = 1 + \lambda$

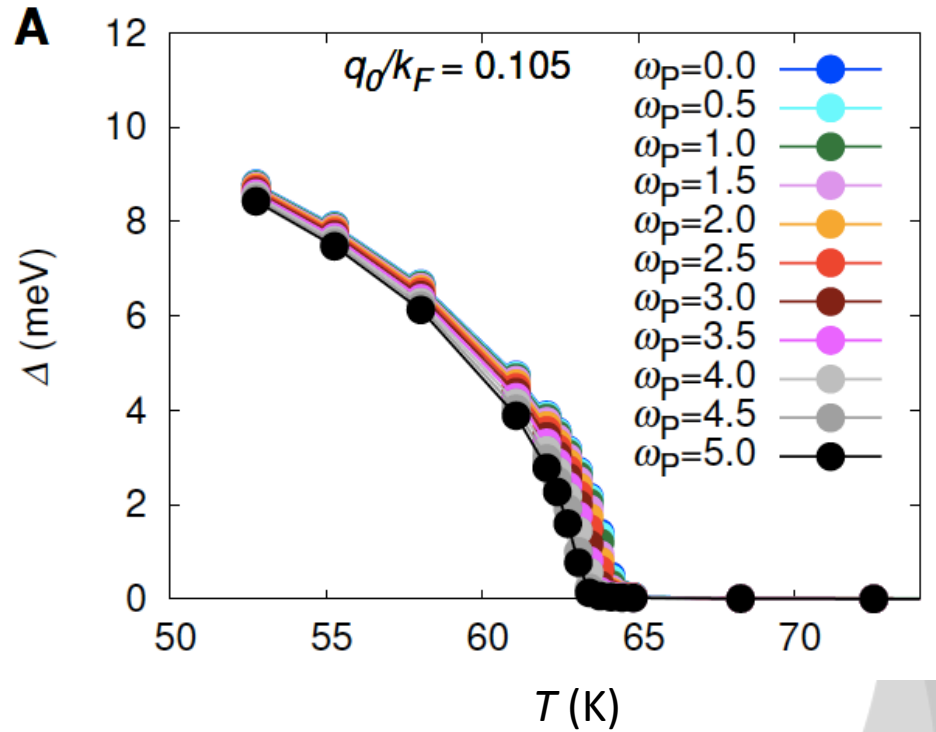
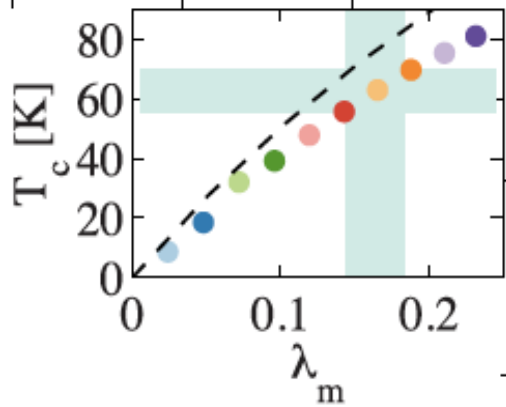
# Cavity materials: Phonon polaritons



enhanced electron-phonon coupling,  
controlled by cavity volume  
( $\omega_p \sim 1/\sqrt{\text{volume}}$ )

# Superconductivity

Rademaker et al., *New J. Phys.* 18, 022001 (2016)



in cavity:  $\lambda$  enhanced, but  $\Omega$  suppressed

suppressed superconductivity despite enhanced el-ph coupling

forward scattering

$$T_C \approx \frac{\lambda \Omega}{2 + 3\lambda}$$

vs.  $T_{C,BCS} \approx 1.13\Omega \exp(-\frac{1}{\lambda})$   
q-independent scattering



# Summary II

- cavity leads to **enhanced electron-phonon coupling**
- can one also enhance superconductivity?

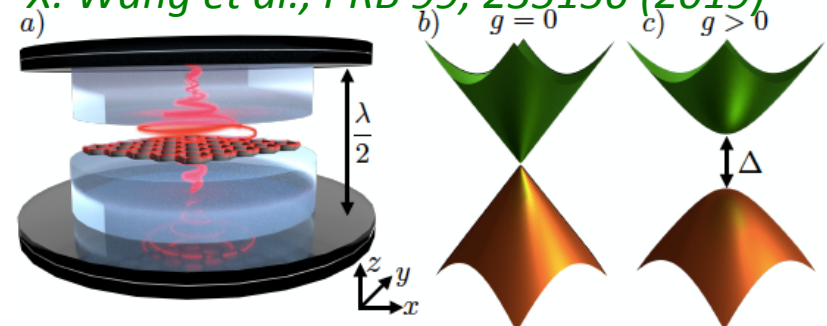
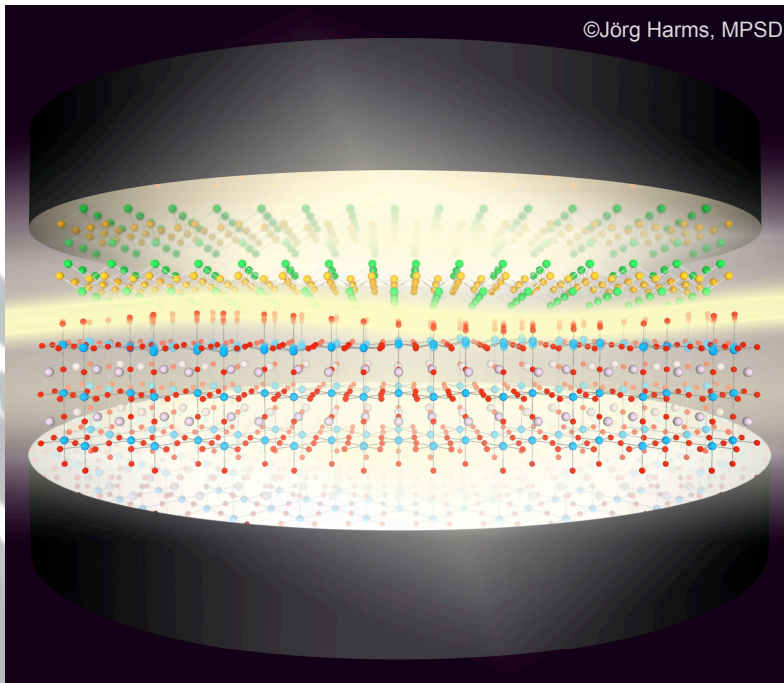
*M. A. Sentef, M. Ruggenthaler, A. Rubio,  
Science Advances 4, eaau6969 (2018)*



M. Ruggenthaler A. Rubio

**Cavity-induced quantum-anomalous  
Hall effect in graphene**

*X. Wang et al., PRB 99, 235156 (2019)*



## manipulation of correlations and topology with light

predicting light-induced  
**Weyl semimetal** in  
pyrochlore iridates

*N. Tancogne-Dejean et al., PRL 121, 097402 (2018)*

*G. E. Topp et al., Nature Comm. 9, 4452 (2018)*

*N. Tancogne-Dejean et al., arXiv:1906.11316*

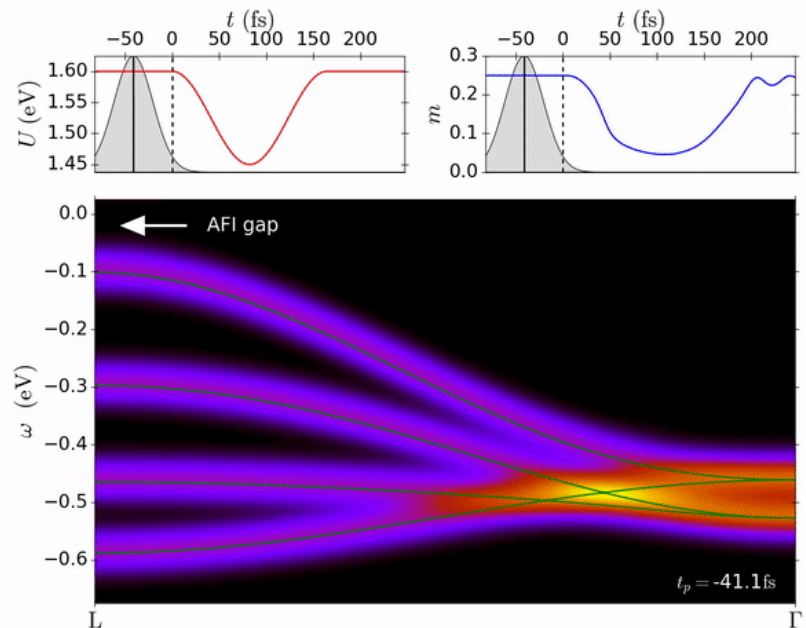


Nicolas Tancogne-Dejean



Gabriel Topp

Max Planck Institute for the Structure and Dynamics of Matter



PERSPECTIVE | TOPOLOGICAL MATTER

### “Weyl”ing away time-reversal symmetry

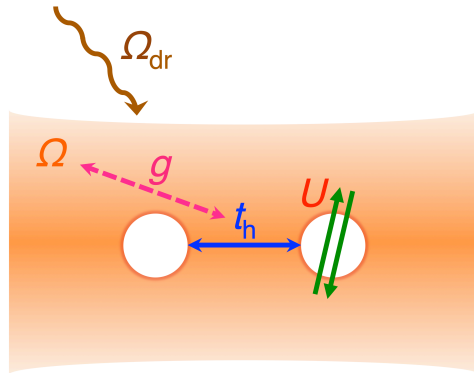
Eduardo H. da Silva Neto

+ See all authors and affiliations

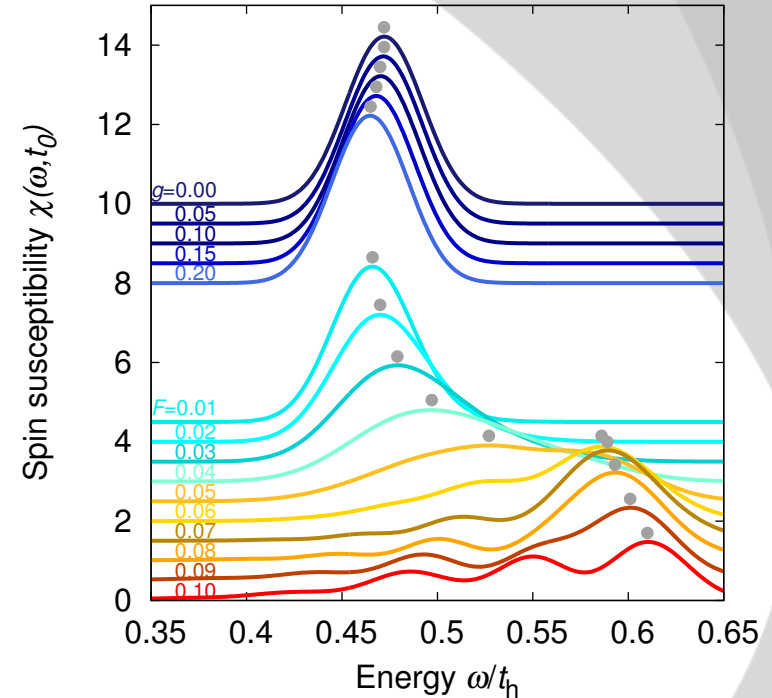
Science 20 Sep 2019:  
Vol. 365, Issue 6459, pp. 1248-1249  
DOI: 10.1126/science.aax6190

also cf. **Gabriel’s talk**

## cavity-Floquet crossover (w/ M. Eckstein's group)



cavity only: cf. Kiffner et al., PRB 2019 (Jacksch group)



manipulate correlations (kinetic exchange coupling)  
with virtual **and** real photons!