

Theory of pump-probe spectroscopy: Ultrafast laser engineering of ordered phases and microscopic couplings

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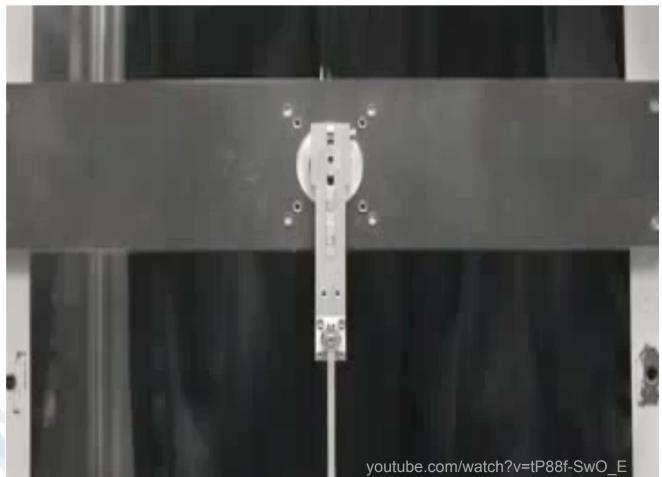


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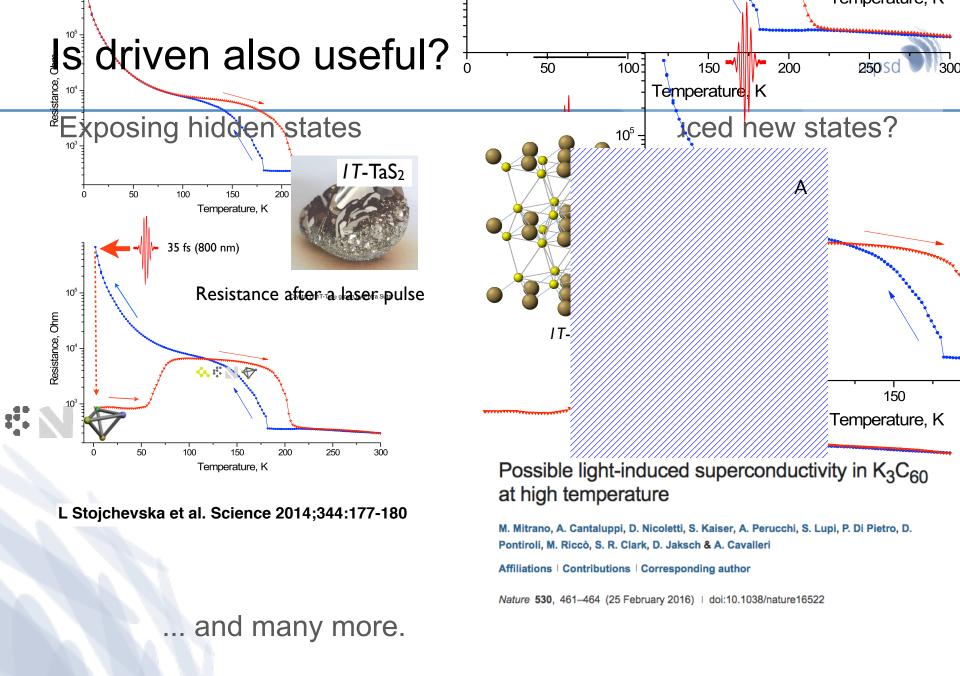
Driven is different



Kapitza pendulum



dynamical stabilization of a metastable state



Ultrafast Materials Science today



Understanding the nature of quasiparticles

- Relaxation dynamics
- Control of couplings

PRB 95, 024304 (2017) PRL 111, 077401 (2013) PRB 95, 205111 (2017) PRX 3, 041033 (2013) PRB 87, 235139 (2013) PRB 90, 075126 (2014) Nat. Comm. 7, 13761 (2016) arXiv:1808.02389

PRL 121, 097402 (2018 arXiv:1802.09437

Understanding or

- Collective oscillat
- Competing order PRB 92, 224517 (2015 1806 08187 arXiv PRB 93, 144506 (2016) arXiv:1808.00712 PRL 118, 087002 (2017) arXiv:1808.04655

100 µm 0.1 cm 10 µm 1 μm Terahertz Infrared Visible Interband Trans. 2DEG Plasmons Phonons Gap phenomena Surface Plasmons 2D & 3D Polarons Carrier Response / Dynamics Dynamics 10^{14} 10¹⁵ THz 10 THz meV 40 meV 0.4eV 4eV 100 fs 10 fs 1 fs ps

Wavelength

Creating new states of matter

Floquet topological states Nature Comm. 6, 7047 (2015) Nature Comm. 8, 13940 (2017) arXiv:1803.07447

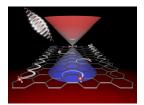


Image courtesy: D. Basov

Outline



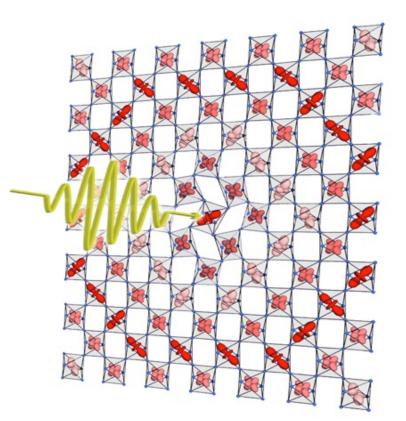
How to modify couplings with light

- Part I: Light-enhanced electron-phonon coupling Resonant excitation of IR phonon enhances electron-phonon coupling E: Pomarico et al., PRB 95, 024304 (2017) – experiment (bilayer graphene) M. A. Sentef, PRB 95, 205111 (2017) – theory
- Part II: Light-reduced Hubbard U

Nonresonant laser driving reduces Hubbard U in NiO N. Tancogne-Dejean et al., PRL 121, 097402 (2018)

I Resonant excitation of crystal lattice

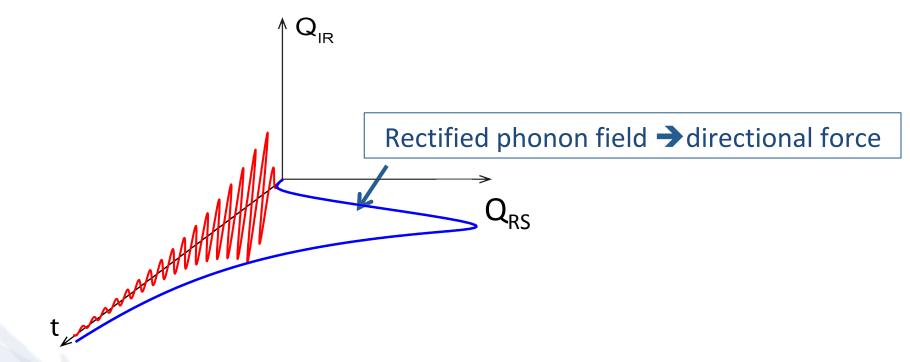




M. Först et al., Nature Physics 7, 854 (2011)

Classical nonlinear phononics





Simplest model: classical dynamics

$$\ddot{Q}_{\rm RS} + \Omega_{\rm RS}^2 Q_{\rm RS} = A Q_{\rm IR}^2$$

$$\ddot{Q}_{\rm IR} + \Omega_{\rm IR}^2 Q_{\rm IR} = \frac{e^* E_0}{\sqrt{M}_{\rm IR}} \sin(\Omega_{\rm IR} t) F(t)$$

"nonlinear phononics"

$$H = AQ_{IR}^2 Q_{RS}$$

M. Först et al., Nature Physics 7, 854 (2011)

Classical nonlinear phononics



Explains a number of observed effects, e.g.,

- structurally induced metal-insulator transitions Rini et al., Nature 449, 72 (2007)
- phononic rectification in YBCO

Mankowsky et al., Nature 516, 71 (2014)

• ferroelectric switching in LiNbO₃

Subedi et al., Phys. Rev. B 89, 220301 (2014)

Mankowsky et al., Phys. Rev. Lett. 118, 197601 (2017)

Classical phonon dynamics **does not** explain all effects in IR-driven materials. examples: - light-induced superconductivity - light-enhanced el-ph coupling ... quantum nature of phonons important? Max Planck Institute for the Structure and Dynamics of Matter

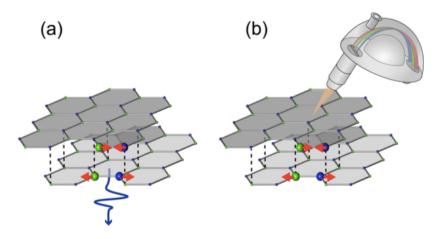
Dynamically enhanced coupling



Enhanced electron-phonon coupling in graphene with periodically distorted lattice

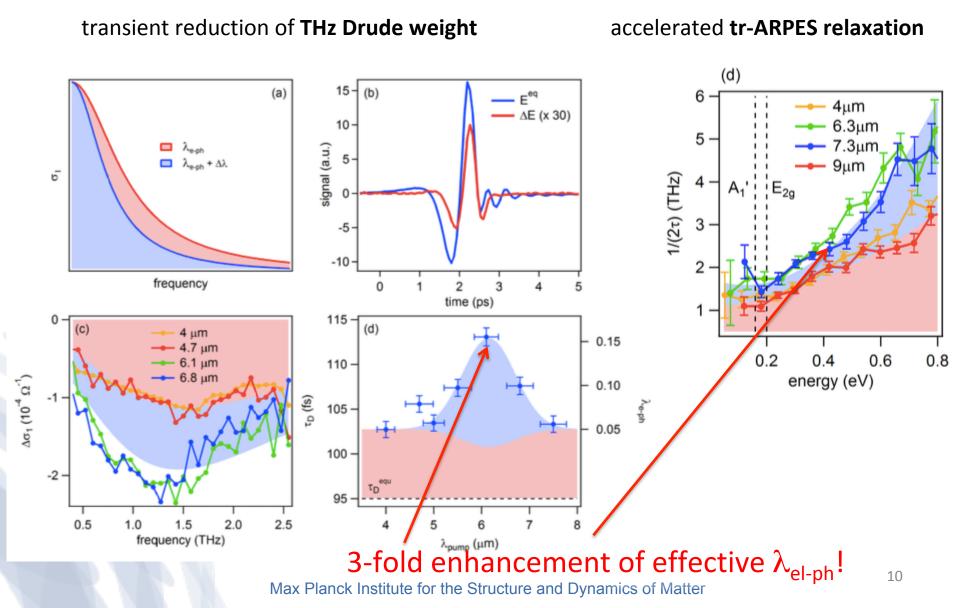
E. Pomarico, M. Mitrano, H. Bromberger, M. A. Sentef, A. Al-Temimy, C. Coletti, A. Stöhr, S. Link, U. Starke, C. Cacho, R. Chapman, E. Springate, A. Cavalleri, and I. Gierz Phys. Rev. B **95**, 024304 – Published 13 January 2017

PRB 95, 024304 (2017) enhanced electron-phonon coupling for pump on resonance with IR phonon



Dynamically enhanced coupling PRB 95, 024304 (2017)





Quantum nonlinear phononics



PRB 95, 205111 (2017)

2-site toy model, solve dynamics exactly

$$\hat{H}(t) = -J \sum_{\sigma} (c_{1,\sigma}^{\dagger} c_{2,\sigma} + c_{2,\sigma}^{\dagger} c_{1,\sigma}) + g_2 \sum_{\sigma,l=1,2}^{\sigma} \hat{n}_{l,\sigma} (b_l + b_l^{\dagger})$$
$$+ \Omega \sum_{l=1,2} b_l^{\dagger} b_l + F(t) \sum_{l=1,2} (b_l + b_l^{\dagger}),$$

also cf. Kennes et al., Nature Physics 13, 479 (2017)

density-dependent squeezing of phonon

Idea: Drive nonlinearly coupled IR-phonon, analyze electronic response

$$\begin{array}{ll} \text{Drive:} & F(t) = F \sin(\omega t), \\ \text{Response:} & I(\omega, t_0) = \operatorname{Re} \int dt_1 \ dt_2 \ e^{i\omega(t_1 - t_2)} s_{t_1, t_2, \tau}(t_0) \\ \text{time-resolved} \\ \text{spectral function} & \times \left[\langle \psi(t_2) | c_{1,\uparrow}^{\dagger} \mathcal{T} e^{-i \int_{t_1}^{t_2} H(t) dt} c_{1,\uparrow} | \psi(t_1) \rangle + \\ & + \langle \psi(t_1) | c_{1,\uparrow} \mathcal{T} e^{-i \int_{t_2}^{t_1} H(t) dt} c_{1,\uparrow}^{\dagger} | \psi(t_2) \rangle \right], \end{array}$$

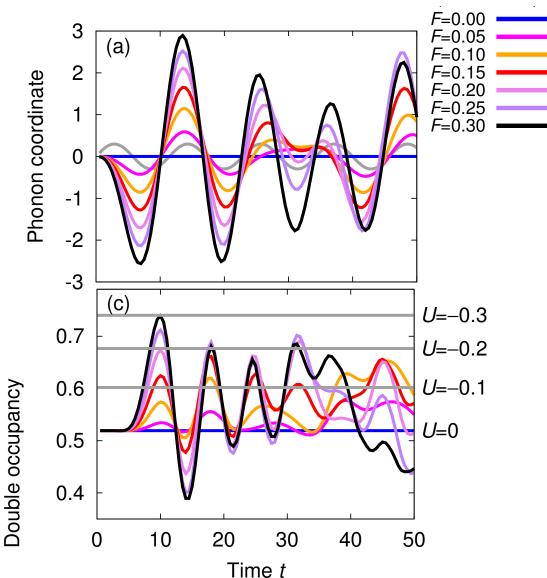
IR-driven nonlinear el-ph system



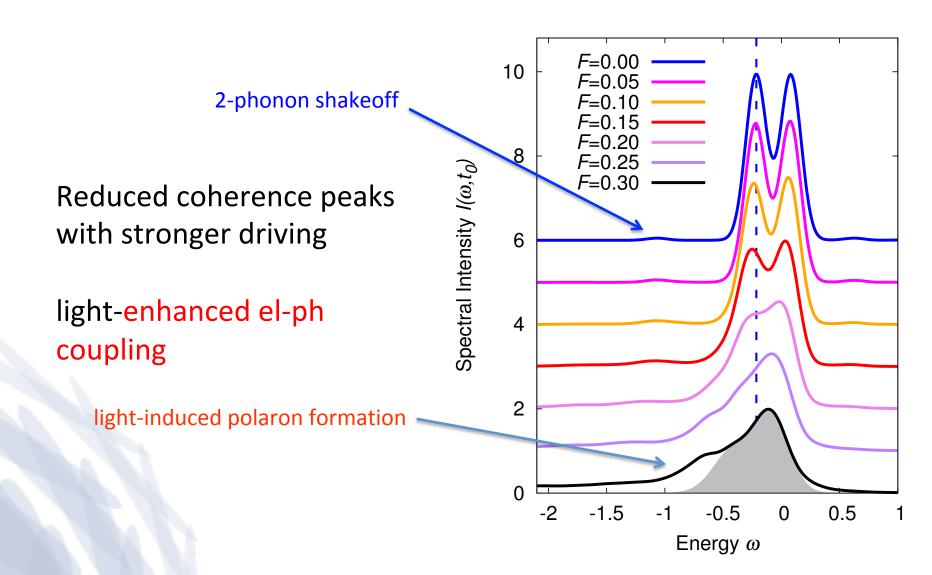
Driving IR phonon with sinusoidal F(t): coherent phonon oscillation

enhancement of local electronic double occupancy

-> induced el-el attraction



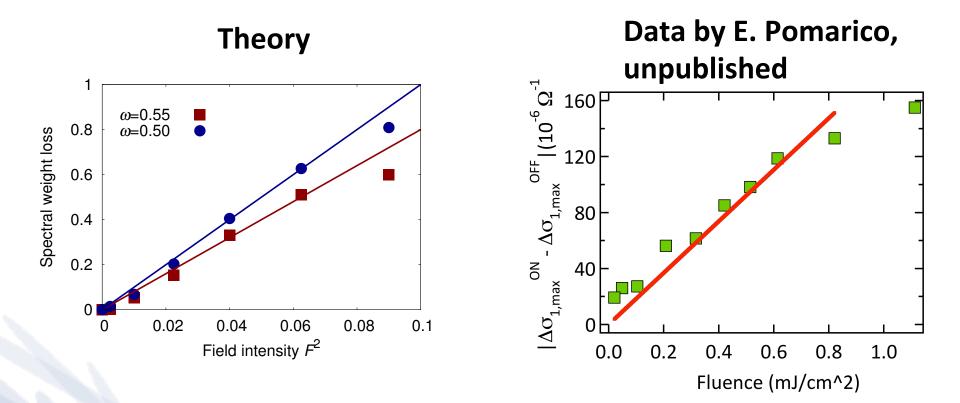
Time-resolved electronic spectrum PRB 95, 205111 (2017) mpsd



Field dependence

PRB 95, 205111 (2017)





Coherence peak weight loss: proportional to field intensity F^2 consistent with experiments



Forced coherent oscillation $\langle \hat{x}_l(t) \rangle \propto F \sin(\omega t)$

q(t)

 $\Sigma =$

Migdal-Eliashberg diagram

effective induced linear coupling

 $\Sigma(t,t') = ig(t)g^*(t')G(t,t')D(t,t')$

time-dependent vertex, amplitude g^2 ~ F^2
=> light-induced coupling, lambda scales ~ F^2





 enhanced electron-phonon coupling in phononically driven bilayer graphene

PRB 95, 024304 (2017)







E. Pomarico

I. Gierz

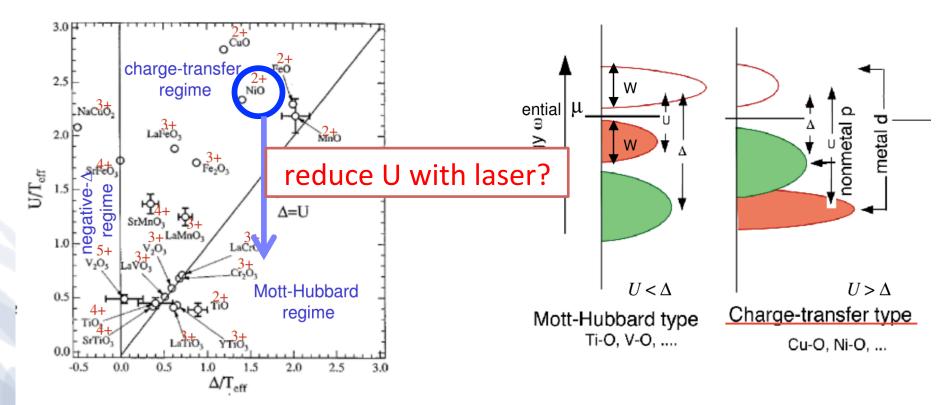
A. Cavalleri

Exact solution of electron-phonon model system:

 theoretical proposal: nonlinear el-ph coupling as mechanism behind this enhancement PRB 95, 205111 (2017)

II Dynamical modification of Hubbard U

Can we drive a charge-transfer insulator towards a Mott insulator?



Zaanen-Sawatzky-Allen phase diagram

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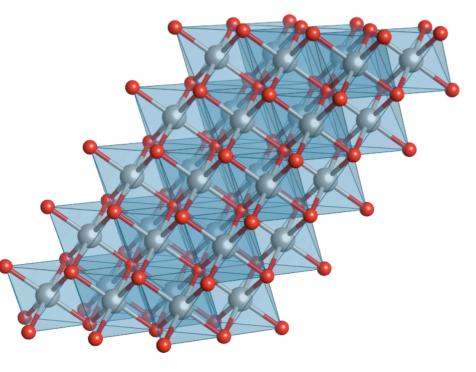
NiO as prototypical charge-transfer insulator mpsd

NiO:

Antiferromagnetic type 2

Band gap: ~4 eV (exp.)

Néel temperature: 523K







DFT with ab initio and self-consistent Hubbard U (Hybrid functional)

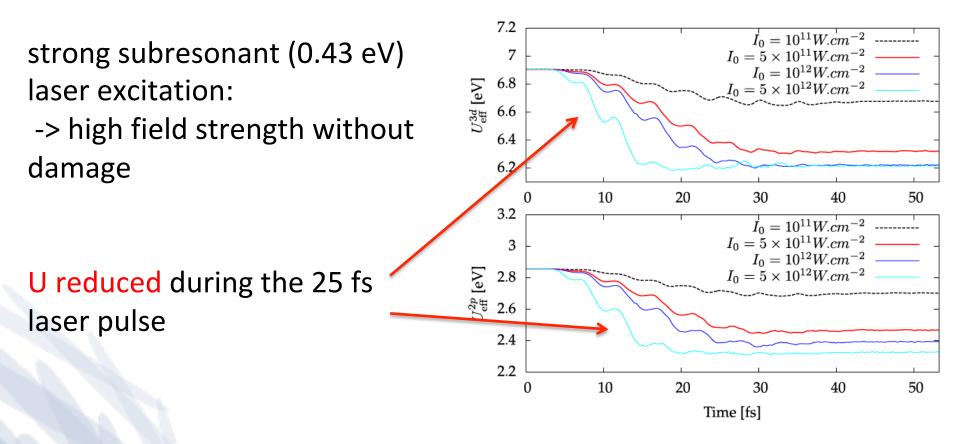
$$E_{\text{DFT}+\text{U}}[n, \{n_{mm'}^{I,\sigma}\}] = E_{\text{DFT}}[n] + E_{ee}[\{n_{mm'}^{I,\sigma}\}] - E_{dc}[\{n_{mm'}^{I,\sigma}\}]$$
Electron-electron interaction Double counting
$$E_{ee} \approx \frac{\bar{U}}{2} \sum_{\{m\},\sigma} N_{m}^{\sigma} N_{m'}^{-\sigma} + \frac{\bar{U} - \bar{J}}{2} \sum_{m \neq m',\sigma} N_{m}^{\sigma} N_{m'}^{\sigma}.$$
Usual expression in DFT+U
$$E_{ee} \approx \frac{\bar{U}}{2} \sum_{\{m\},\sigma} N_{mm'}^{\sigma} \bar{P}_{mm'}^{\sigma} \bar{P}_{m'm''}^{\sigma} (mm'|m''m'') - \frac{1}{2} \sum_{\{m\}} \sum_{\alpha} \frac{\bar{P}_{mm'}^{\alpha} \bar{P}_{m''m''}^{\alpha} (mm''|m''m')}{\bar{P}_{mm'}^{\alpha} \bar{P}_{m''m'''}^{\alpha} (mm''|m''m')}$$

$$ACBNO \text{ functional}$$

PRX 5,011006 (2015)

- alternative to constrained RPA
- numerically efficient
- direct extension to time-dependent case (adiabatic approximation)

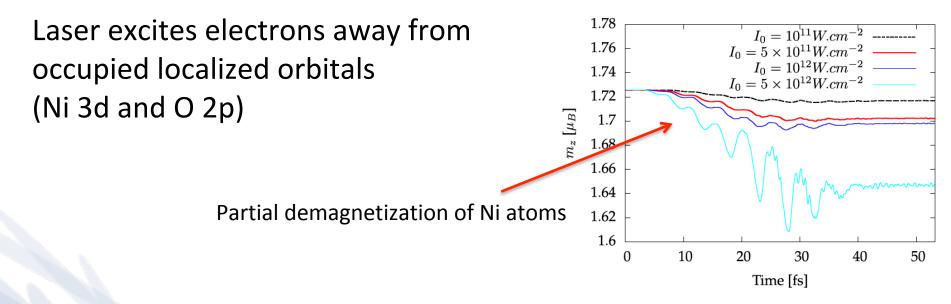




Typical intensities in strong field physics in solids



U measures the Coulomb interaction screened by itinerant electrons



- Light-enhanced screening
- Decrease of U





 Ultrafast reduction of Hubbard U in NiO via induced extra screening

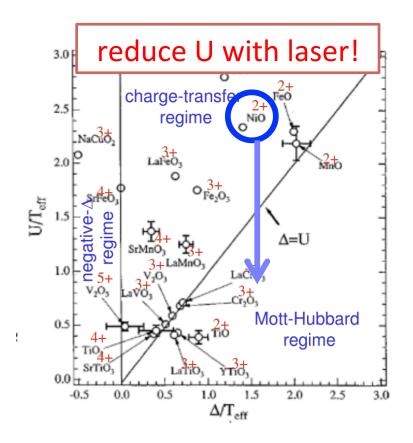
N. Tancogne-Dejean et al., PRL 121, 097402 (2018)



N. Tancogne-Dejean

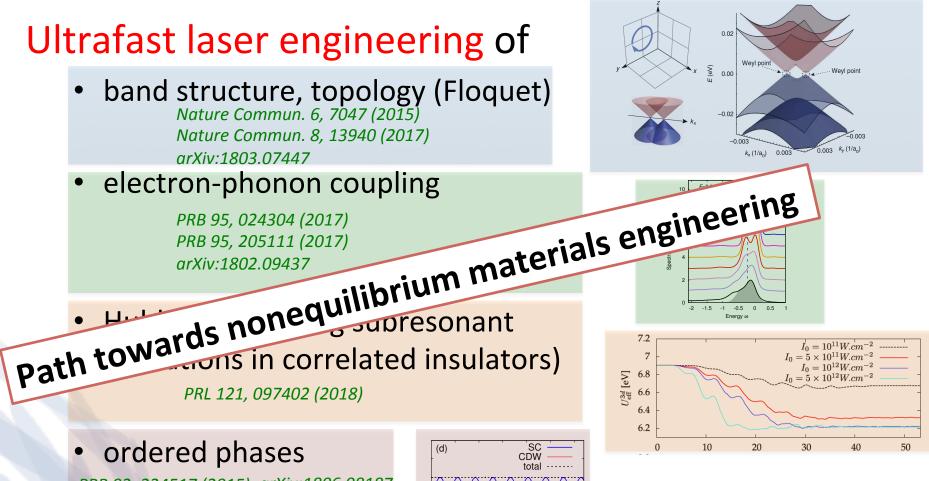


A. Rubio



Summary





PRB 92, 224517 (2015) arXiv:1806.08187 PRB 93, 144506 (2016) arXiv:1808.00712 PRL 118, 087002 (2017) arXiv:1808.04655

