

### Theory of pump-probe spectroscopy: Ultrafast laser engineering of ordered phases and microscopic couplings

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## Pump-probe spectroscopy (1887)



stroboscopic investigations of dynamic phenomena





Muybridge 1887

## Pump-probe spectroscopy (today)

- stroboscopic investigations of dynamic phenomena



Simulations of time-resolved ARPES: PRX 3, 041033 (2013), PRB 90, 075126 (2014), PRB 92, 224517 (2015), Nature Commun. 7, 13761 (2016)

Image courtesy: J. Sobota / F. Schmitt

## Ultrafast Materials Science today

#### Understanding the nature of quasiparticles

- Relaxation dynamics
- Control of couplings

PRL 111, 077401 (2013)PRB 95, 024304 (2017)PRX 3, 041033 (2013)PRB 95, 205111 (2017)PRB 87, 235139 (2013)arXiv:1712.01067PRB 90, 075126 (2014)arXiv:1802.09437Nature Commun. 7, 13761 (2016)

#### Understanding or

- Collective oscillat
- Competing order *PRB 92, 224517 (2015 <i>PRB 93, 144506 (2016) PRL 118, 087002 (2017)*

#### Creating new states of matter

Floquet topological states

Nature Commun. 6, 7047 (2015) Nature Commun. 8, 13940 (2017)



Image courtesy: D. Basov





## Outline





## Outline



## How to modify couplings with light

- Part I: Light-enhanced electron-phonon coupling Resonant excitation of IR phonon enhances electron-phonon coupling E: Pomarico et al., PRB 95, 024304 (2017) – experiment (bilayer graphene) M. A. Sentef, PRB 95, 205111 (2017) – theory
- Part II: Light-reduced Hubbard U

Nonresonant laser driving reduces Hubbard U in NiO

N. Tancogne-Dejean et al., 1712.01067

## I Resonant excitation of crystal lattice





M. Först et al., Nature Physics 7, 854 (2011)

## **Classical nonlinear phononics**





Simplest model: classical dynamics

$$\ddot{Q}_{\rm RS} + \Omega_{\rm RS}^2 Q_{\rm RS} = A Q_{\rm IR}^2$$

$$\ddot{Q}_{\rm IR} + \Omega_{\rm IR}^2 Q_{\rm IR} = \frac{e^* E_0}{\sqrt{M}_{\rm IR}} \sin(\Omega_{\rm IR} t) F(t)$$

"nonlinear phononics"

$$H = AQ_{IR}^2 Q_{RS}$$

M. Först et al., Nature Physics 7, 854 (2011)

## **Classical nonlinear phononics**



Explains a number of observed effects, e.g.,

- structurally induced metal-insulator transitions Rini et al., Nature 449, 72 (2007)
- phononic rectification in YBCO

Mankowsky et al., Nature 516, 71 (2014)

• ferroelectric switching in LiNbO<sub>3</sub>

Subedi et al., Phys. Rev. B 89, 220301 (2014)

Mankowsky et al., Phys. Rev. Lett. 118, 197601 (2017)

Classical mechanistic phonon dynamics does not explain all effects in IR-driven materials. examples: - light-induced superconductivity - light-enhanced el-ph coupling ... quantum nature of phonons important? Light-induced superconductivity?



M. Mitrano et al., Nature 530, 461 (2016)

Lattice control of reflectivity in  $K_3C_{60}$ 

<sup>10</sup> Not (easily) explained by classical nonlinear phononics

#### Hard problem!

Kennes, Millis, Knap, Demler, Murakami, Eckstein,

Werner, Thorwart, Mazza, Georges, Fabrizio, Galitskii, Sentef, Kollath, ... 0.0

#### Simpler question: what happens to electron-phonon coupling under IR driving in a metal?

## Dynamically enhanced coupling



Enhanced electron-phonon coupling in graphene with periodically distorted lattice

E. Pomarico, M. Mitrano, H. Bromberger, M. A. Sentef, A. Al-Temimy, C. Coletti, A. Stöhr, S. Link, U. Starke, C. Cacho, R. Chapman, E. Springate, A. Cavalleri, and I. Gierz Phys. Rev. B **95**, 024304 – Published 13 January 2017

#### *PRB 95, 024304 (2017)* enhanced electron-phonon for pump on resonance with IR phonon



# Dynamically enhanced coupling PRB 95, 024304 (2017)





## Quantum nonlinear phononics



PRB 95, 205111 (2017)

2-site toy model, solve dynamics exactly

$$\begin{split} \hat{H}(t) &= -J \sum_{\sigma} (c_{1,\sigma}^{\dagger} c_{2,\sigma} + c_{2,\sigma}^{\dagger} c_{1,\sigma}) \\ &+ g_2 \sum_{\sigma,l=1,2}^{\sigma} \hat{n}_{l,\sigma} (b_l + b_l^{\dagger}) \\ &+ \Omega \sum_{l=1,2} b_l^{\dagger} b_l + F(t) \sum_{l=1,2} (b_l + b_l^{\dagger}) \end{split}$$

also cf. Kennes et al., Nature Physics 13, 479 (2017)

electron-occupation dependent squeezing of phonon;

*g*<sub>2</sub> can be positive or negative in materials -> mode hardening or softening

Idea: Drive nonlinearly coupled IR-phonon, analyze electronic response

$$\begin{array}{ll} \text{Drive:} & F(t) = F \sin(\omega t), \\ \text{Response:} & I(\omega, t_0) = \operatorname{Re} \int dt_1 \ dt_2 \ e^{i\omega(t_1 - t_2)} s_{t_1, t_2, \tau}(t_0) \\ \text{time-resolved} \\ \text{spectral function} & \times \left[ \langle \psi(t_2) | c_{1,\uparrow}^{\dagger} \mathcal{T} e^{-i \int_{t_1}^{t_2} H(t) dt} c_{1,\uparrow} | \psi(t_1) \rangle + \\ & + \langle \psi(t_1) | c_{1,\uparrow} \mathcal{T} e^{-i \int_{t_2}^{t_1} H(t) dt} c_{1,\uparrow}^{\dagger} | \psi(t_2) \rangle \right], \end{array}$$

## IR-driven nonlinear el-ph system



Driving IR phonon with sinusoidal F(t): coherent phonon oscillation

enhancement of local electronic double occupancy

-> induced el-el attraction



Time-resolved electronic spectrum PRB 95, 205111 (2017) mpsd



## Field scaling

PRB 95, 205111 (2017)





## Coherence peak weight loss: proportional to field intensity F^2 consistent with experiments



Forced coherent oscillation  $\langle \hat{x}_l(t) \rangle \propto F \sin(\omega t)$ 

q(t)

 $\Sigma =$ 

Migdal-Eliashberg diagram

effective induced linear coupling

 $\Sigma(t,t') = ig(t)g^*(t')G(t,t')D(t,t')$ 

time-dependent vertex, amplitude g^2 ~ F^2
=> light-induced coupling, lambda scales ~ F^2





 enhanced electron-phonon coupling in phononically driven bilayer graphene

PRB 95, 024304 (2017)





E. Pomarico

I. Gierz

A. Cavalleri

Exact solution of electron-phonon model system:

 theoretical proposal: nonlinear el-ph coupling as mechanism behind this enhancement PRB 95, 205111 (2017)

#### Cavity QED superconductivity arXiv:1802.09437





Materials engineering in nanocavities through coupling to quantum light



M. Ruggenthaler A. Rubio





II Dynamical modification of Hubbard U

# Can we drive a charge-transfer insulator towards a Mott insulator?



Zaanen-Sawatzky-Allen phase diagram

Max Planck Institute for the Structure and Dynamics of Matter

## NiO as prototypical charge-transfer insulator mpsd

NiO:

Antiferromagnetic type 2

Band gap: ~4 eV (exp.)

Néel temperature: 523K







DFT with ab initio and self-consistent Hubbard U (Hybrid functional)

$$E_{\text{DFT+U}}[n, \{n_{mm'}^{I,\sigma}\}] = E_{\text{DFT}}[n] + E_{ee}[\{n_{mm'}^{I,\sigma}\}] - E_{dc}[\{n_{mm'}^{I,\sigma}\}]$$
Electron-electron interaction Double counting
$$E_{ee} \approx \frac{\bar{U}}{2} \sum_{\{m\},\sigma} N_m^{\sigma} N_{m'}^{-\sigma} + \frac{\bar{U} - \bar{J}}{2} \sum_{m \neq m',\sigma} N_m^{\sigma} N_{m'}^{\sigma}.$$
Usual expression in DFT+U
$$E_{ee} \approx \frac{\bar{U}}{2} \sum_{\{m\},\sigma} N_m^{\sigma} N_{m'}^{-\sigma} + \frac{\bar{U} - \bar{J}}{2} \sum_{m \neq m',\sigma} N_m^{\sigma} N_{m'}^{\sigma}.$$

$$E_{ee} \approx \frac{1}{2} \sum_{\{m\}} \sum_{\alpha,\beta} \frac{\bar{P}_{mm'}^{\alpha} \bar{P}_{m''m''}^{\beta} (mm'|m''m'')}{-\frac{1}{2} \sum_{\{m\}}} \sum_{\alpha,\beta} \frac{\bar{P}_{mm'}^{\alpha} \bar{P}_{m''m''}^{\beta} (mm''|m''m'')}{\bar{P}_{mm'}^{\alpha} \bar{P}_{m''m''}^{\alpha} (mm''|m''m'')}$$

$$ACBNO \text{ functional PRX 5,011006 (2015)}$$

alternative to constrained RPA

- numerically efficient
- direct extension to time-dependent case (adiabatic approximation)





Typical intensities in strong field physics in solids



U measures the Coulomb interaction screened by itinerant electrons



- Polarization of itinerant electrons increases
- Enhanced screening
- Decrease of U

## Summary II



- Ultrafast reduction of Hubbard U in NiO via induced extra screening
- Towards light-induced Mott insulators?
- N. Tancogne-Dejean et al., 1712.01067



N. Tancogne-Dejean



A. Rubio



#### Laser-controlled competing orders Phys. Rev. Lett. 118, 087002 (2017)





Controlling competing orders by driving near gap resonance



# Nonthermal magnetic Weyl semimetal







Nonthermal pathway to magnetic Weyl semimetal in pyrochlore iridates



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... and many more



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## Summary



